Thermal properties of single-phase Y$_2$SiO$_5$

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Abstract

Y$_2$SiO$_5$ is a promising candidate for oxidation-resistant or environmental/thermal barrier coatings (ETBC) due to its excellent high-temperature stability, low elastic modulus and low oxygen permeability. In this paper, we investigated the thermal properties of Y$_2$SiO$_5$ comprehensively, including thermal expansion, thermal diffusivity, heat capacity and thermal conductivity. It is interesting that Y$_2$SiO$_5$ has a very low thermal conductivity ($\sim 1.40$ W/m K) but a relatively high linear thermal expansion coefficient ($8.36 \pm 0.5 \times 10^{-6}$ K$^{-1}$), suggesting compatible thermal and mechanical properties to some non-oxide ceramics and nickel superalloys as ETBC layer. Y$_2$SiO$_5$ is also an ideal EBC on YSZ TBC layer due to their close thermal expansion coefficients. As a continuous source of Y$^{3+}$, it is predicted that Y$_2$SiO$_5$ EBC may prolong the lifetime of zirconia-based TBC by stopping the degradation aroused by the loss of Y stabilizer.

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Keywords: Y$_2$SiO$_5$; Thermal properties; Silicate; Thermal conductivity; Engine components

1. Introduction

Si-based ceramics, such as silicon carbide and silicon nitride, have superior high-temperature strength and creep resistance, and show potential to improve gas turbine engine performance. However, the lack of environmental durability in combustion environments impedes the applications of these materials. Therefore, a protective coating is in need. Recently, experimental trials have proven that Y$_2$Si$_2$O$_7$, Y$_2$SiO$_5$ as well as Lu$_2$Si$_2$O$_7$, Yb$_2$Si$_2$O$_7$ are promising candidates for environmental barrier coatings (EBC) on silicon-based ceramics. Yttrium silicate (Y$_2$SiO$_5$) is an important laser crystal that has been synthesized by a chemical method since 1963. Most research activities are related to rare earth (Ce$^{3+}$, Eu$^{3+}$, etc.) doped Y$_2$SiO$_5$ crystals to be used as blue phosphor or Cr$^{4+}$-doped Y$_2$SiO$_5$ as a saturable-absorber Q-switch laser. As an important phase in the SiO$_2$–Y$_2$O$_3$–Si$_3$N$_4$ phase diagram, Y$_2$SiO$_5$ is frequently identified as a precipitated phase at the grain boundaries of sintered Si$_3$N$_4$ with Y$_2$O$_3$ or Y$_2$O$_3$ + SiO$_2$ as sintering aids. And it is confirmed that Y$_2$SiO$_5$ can improve the high-temperature performance of silicon nitride. The melting point of Y$_2$SiO$_5$ is up to 1950$^\circ$C guaranteeing its potential high-temperature structural applications. Although the spectroscopic studies of doped Y$_2$SiO$_5$ have been widely performed, the fundamental properties of the host crystal itself are still in dearth.

Recently, Y$_2$SiO$_5$ has attracted wide attentions and a number of published works have focused on its coating applications, which were aroused by two interesting characteristics of Y$_2$SiO$_5$. Firstly, Y$_2$SiO$_5$ has a low oxygen permeability over a wide temperature range, as reported by Ogura et al., the oxygen permeability constant is $10^{-10}$ kg/(m s) at 773 K. Secondly, the thermal expansion coefficient of Y$_2$SiO$_5$ (Table 1) matches well with some non-oxide materials such as silicon-based ceramics, C/C composites and so on, especially in combination with Y$_2$Si$_2$O$_7$. However, the preparation of single-phase and fully dense Y$_2$SiO$_5$ bulk material is very difficult. The synthesis of Y$_2$SiO$_5$ is always accompanied with the presence of undesired compounds such as Y$_{4.67}$(SiO$_4$)$_3$O, Y$_2$Si$_2$O$_7$, etc., as well as incomplete polymorph transformation between the low-temperature phase (X1–Y$_2$SiO$_5$) and high-temperature phase (X2–Y$_2$SiO$_5$). Impurities have a strong influence on the final properties of Y$_2$SiO$_5$. For example, the measured thermal properties of single-phase Y$_2$SiO$_5$.
Table 1
Reported thermal expansion coefficients for Y$_2$SiO$_5$

<table>
<thead>
<tr>
<th>Material</th>
<th>Temperature (K)</th>
<th>TEC ($\times 10^{-6}$ K$^{-1}$)</th>
<th>Reference</th>
</tr>
</thead>
<tbody>
<tr>
<td>Single crystal Y$_2$SiO$_5$</td>
<td>300–1000</td>
<td>$\alpha_a = 0.6$, $\alpha_b = 7.5$, $\alpha_c = 11.4$</td>
<td>O'Bryan et al.$^{17}$</td>
</tr>
<tr>
<td></td>
<td>293–1123</td>
<td>$\Delta V/\Delta T = 17$</td>
<td></td>
</tr>
<tr>
<td>Y$_2$SiO$_5$ powders</td>
<td>293–1123</td>
<td>$\alpha_a = 5.2$, $\alpha_b = 7.5$, $\alpha_c = 10.7$</td>
<td></td>
</tr>
<tr>
<td></td>
<td>1123–1673</td>
<td>$\alpha_a = 19.5$, $\alpha_b = 12.5$, $\alpha_c = 25$</td>
<td>Nowok et al.$^{18}$</td>
</tr>
<tr>
<td>X1-Y$_2$SiO$_5$ powders</td>
<td>394–1273</td>
<td>$\alpha = 5.0 \times 0.00421T$</td>
<td>Fukuda and Matsubara$^{19}$</td>
</tr>
<tr>
<td>X2-Y$_2$SiO$_5$ powders</td>
<td>394–1473</td>
<td>$\alpha = 5.0 \times 0.00257T$</td>
<td></td>
</tr>
<tr>
<td>Polycrystalline bulk material</td>
<td>300–1623</td>
<td>$\alpha = 5.3$</td>
<td>Ogura et al.$^{20}$</td>
</tr>
<tr>
<td>(Y$_2$SiO$_5$ + Y$_3$Si$_2$O$_7$)</td>
<td>300–1273</td>
<td>$\alpha = 6.9$</td>
<td>Apricio and Duran$^{21}$</td>
</tr>
<tr>
<td>Polycrystalline bulk material</td>
<td>300–1173</td>
<td>$\alpha = 3.05$</td>
<td>Wagner et al.$^{22}$</td>
</tr>
<tr>
<td>(Y$_2$SiO$_5$ + Y$_2$O$_3$)</td>
<td>1173–1673</td>
<td>$\alpha = 4.09$</td>
<td></td>
</tr>
<tr>
<td>Single-phase polycrystalline</td>
<td>300–1573</td>
<td>$\alpha = 8.36$</td>
<td>Present work</td>
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<tr>
<td>Y$_2$SiO$_5$ bulk material</td>
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</tr>
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</table>

expansion coefficients of Y$_2$SiO$_5$ in Table 1 show a wide deviation from author to author, which was mainly caused by the undesired impurities. The fundamental properties such as mechanical and thermal properties for single-phase Y$_2$SiO$_5$ are seldom investigated. But these data are indispensable for selection of Y$_2$SiO$_5$ as candidate for both high-temperature structural components and coatings.

In our previous paper, we successfully synthesized single-phase Y$_2$SiO$_5$ powders (X2-Y$_2$SiO$_5$, hereafter written as Y$_2$SiO$_5$ for brevity) from Y$_2$O$_3$ and SiO$_2$ powders,$^{25}$ and evaluated its mechanical properties.$^{26}$ Y$_2$SiO$_5$ has low Young’s modulus but high damage tolerance, which endow the promising coating applications of Y$_2$SiO$_5$ from the view of mechanical properties. In this paper, the thermal expansion coefficient of polycrystalline Y$_2$SiO$_5$ was investigated by measuring the dimension variation of cylindrical specimens using push-rod thermal expansion dilatometer. Other thermal properties, e.g., thermal diffusivity, thermal capacity and thermal conductivity were also studied. These thermal property data are helpful for the further design of Y$_2$SiO$_5$ as coatings or high-temperature structural components.

2. Experimental procedure

Y$_2$SiO$_5$ powders were synthesized via a solid/liquid reaction method. The initial yttria and silica powders were calcined with 3 mol.% LiYO$_2$ additive at 1500 $^\circ$C for 2 h in air. Bulk material with 98% theoretical density was obtained by pressureless sintering of the synthesized Y$_2$SiO$_5$ powders at 1500 $^\circ$C for 60 min in air. Details of the preparation can be found in our previous work.$^{25}$ Fig. 1(a) presents the microstructure of the thermal etched surface of Y$_2$SiO$_5$ material via heating to 1300 $^\circ$C with a rate of 10 $^\circ$C/min, holding for 30 min, and then cooling to room temperature with a rate of 15–20 $^\circ$C/min. The sample displays equiaxed grains in the majority with an average grain size of

Fig. 1. (a) Microstructure of the thermal etched surface of Y$_2$SiO$_5$, and (b) XRD pattern collected on the sintered material.
about 6 µm. A few residual small pores can be found in the sintered material, and all the pores exist at the grain boundaries. Fig. 1(b) shows a XRD pattern collected on the sample after sintering. It is clear that the sintered material is single-phase Y$_2$SiO$_5$ and without visible preferred orientation.

The average linear thermal expansion coefficient was obtained from the temperature dependent changes of the length of the specimen from room temperature to 1573 K in air by using a high-temperature dilatometer (Setaram Setsys 24, Caluire, France). The data were continuously recorded during heating at a rate of 2 K/min as well as cooling processes. The specimens were cylindrical pieces with 6 mm in diameter and 8 mm in height.

The thermal diffusivity measurement was conducted on a Flashline 5000 thermal properties analyzer (Anter Corporation, Pittsburg, USA) according to the procedure described in ASTM E1461-01, which is based on delivering a pulse of thermal energy to one face of the analyzed sample and monitoring the temperature rise of the opposite face. The thermal diffusivity, $D_{th}$ (cm$^2$/s) is evaluated from sample thickness $L$ (cm), and the time, $t_{1/2}$ (s), elapsed from the moment when the pulse of thermal energy is delivered to the moment when the temperature of the opposite face of the sample reaches half of its maximum:

$$D_{th} = \frac{0.13879L^2}{t_{1/2}}$$

(1)

To measure the constant pressure molar heat capacity $C_p$ and thermal conductivity $\kappa$, a disc sample (Ø 12.7 mm × 1.0 mm) was used. The thermal conductivity was determined in the temperature range from 473 to 1273 K. Prior to the thermal diffusion test the samples were sprayed with a thin layer of colloidal graphite to ensure complete and uniform absorption of the laser pulse. Three measurements were taken at each temperature and the data were recorded using an Anter FL5000 software. Employing a multi-sample configuration system and testing a reference sample (graphite) adjacent to Y$_2$SiO$_5$, the heat capacity of Y$_2$SiO$_5$ can be obtained in parallel with thermal diffusivity. Then, the thermal diffusivity results were converted to thermal conductivities using the heat capacity results and measured density of Y$_2$SiO$_5$. Thermal conductivity was calculated from the values of thermal diffusivity, specific heat capacity, and measured density, according to the relation:

$$\kappa(T) = D_{th}(T)C_p(T)\rho(T)$$

(2)

3. Results and discussions

3.1. Thermal properties of Y$_2$SiO$_5$

Thermal expansion of polycrystalline Y$_2$SiO$_5$ measured by the push-rod dilatometer was illustrated in Fig. 2. With the increasing of temperature to 1573 K, the length of Y$_2$SiO$_5$ samples expanded linearly. Since the as-sintered sample has no preferable orientation (Fig. 1(b)), an isotropic thermal expansion for Y$_2$SiO$_5$ bulk materials was observed. The fractional change in the linear thermal expansion of Y$_2$SiO$_5$, $\Delta L/L_{300\,K} = (L - L_{300\,K})/L_{300\,K}$, expressed as a percent is shown in Fig. 2. The linear thermal expansion coefficient was determined in terms of the slope of thermal expansion vs. temperature during heating and cooling, and the linear CTE for γ-Y$_2$Si$_2$O$_7$ is $(8.36 \pm 0.5) \times 10^{-6}$ K$^{-1}$. This value is lower than that obtained from the high-temperature power XRD measurement (Nowok et al.) but slightly higher than those measured on the polycrystalline samples containing impurities like Y$_2$Si$_2$O$_7$ or Y$_{2.86}$(SiO$_4$)$_3$O which has low TEC value (Table 1). Moreover, no abnormal expansion was observed in Fig. 2 in the examined temperature range, though a phase transformation from the low-temperature polymorph X1-Y$_2$SiO$_5$ to the high-temperature X2-Y$_2$SiO$_5$ around 1173 K was reported in some literatures.

Fig. 3 is the thermal diffusivity data fitted by the second order polynomial equation, and the obtained result takes on a form like:

$$D_{th} = 0.01124 - 1.314 \times 10^{-5}T + 6.331 \times 10^{-9}T^2$$

(3)

with a reliability value of $R^2 = 0.998$.

The temperature dependence of the heat capacity of Y$_2$SiO$_5$ is plotted in Fig. 4, in which the data from Y$_2$O$_3$ and SiO$_2$ are also given for comparison. Curve fitting of the experimental data

![Fig. 2. Linear thermal expansion of polycrystalline Y$_2$SiO$_5$ bulk material.](image)

![Fig. 3. Thermal diffusivity versus temperature for Y$_2$SiO$_5$.](image)
of $Y_2SiO_5$ yields:

$$C_p = 177.48 + 29.68 \times 10^{-3} T - 59.2 \times 10^5 T^{-2}$$

with a $R^2$ value of 0.994. The molar heat capacities extrapolated from this relationship are 120 and 215 J/(mol K) at room temperature (300 K) and high temperature (1400 K), respectively. The obtained heat capacity data of $Y_2SiO_5$ coincides will with the summation of $Y_2O_3 + SiO_2$ except for a minor deviation at the low and high limiting temperatures. Based on the Neumann–Kopp law, the heat capacity of $Y_2SiO_5$ can be determined as a summation of the atomic heat capacities of the constituent elements.30 At temperatures above the Debye temperature (the Debye temperature for $Y_2SiO_5$ is 580 K, the calculation will be mentioned in later section), the molar heat capacity at constant volume, $C_v$, of each atom is about 3R ($\approx 25 J/(mol K)$) as predicted by the Dulong–Petit equation, where $R$ is the gas constant.31 Here, the calculated temperature-independent heat capacity of $Y_2SiO_5$ above Debye temperature is approximately 200 J/(mol K), which is slightly lower than the measured value at high temperatures (215 J/(mol K)) without taking into account of the difference between $C_p$ and $C_v$ (at room temperature, $(C_p - C_v)$ for a typical solid is about 1–8 J/(mol K) and it reaches a maximum at the melting point with a value of about 10% of $C_v$).30

The thermal conductivity of $Y_2SiO_5$ is plotted in Fig. 5. An inverse proportional fitting of the data yields the following relationship:

$$\kappa = 1.138 + \frac{216.2}{T}$$

From this equation, the thermal conductivities of $Y_2SiO_5$ are almost inverse proportional to temperature and are extrapolated to be 1.86 and 1.29 W/m K, respectively, at 300 K and 1400 K. As the microstructure and XRD of $Y_2SiO_5$ presented in Fig. 1, only a few residual small pores were observed at the grain boundaries and no grain-boundary phase was detected either by SEM or by XRD. The quantity of grain-boundary pores is too small to influence the thermal conductivity, and thus the measured thermal conductivity should be the intrinsic value of $Y_2SiO_5$. Therefore, $Y_2SiO_5$ has a thermal conductivity much lower than most commonly used thermal barrier coating materials, such as YSZ, La$_2$Zr$_2$O$_7$, etc.32,33 Fig. 6 exhibits the thermal conductivities of some selected TBC materials as a function of temperatures.5,32–34 Though a dense $Y_2SiO_5$ material was used for thermal conductivity measurement, its conductivity is only slightly higher than those of APS 7YSZ (air plasma sprayed 7 mol.% yttria stabilized zirconia) coating and monoclinic zirconia.32

### 3.2. Minimum high-temperature thermal conductivity

The thermal conductivity of a material describes the diffusivity of heat flow by phonon transport on temperature gradient. In the first successful model for thermal conductivity, Debye used an analogy of the kinetic theory of gases to derive an expression of the thermal conductivity as:

$$\kappa = \frac{C_v v_m A}{3}$$

Fig. 6. Thermal conductivity of some selected thermal barrier coating candidates.
where \( C_v \) is the specific heat, \( v_m \) the speed of sound and \( \Lambda \) the phonon mean free path. In the subsequent developments, a temperature that is characteristic of the material and now known as the Debye temperature, \( \theta_D \), plays a central role in phonon theories of thermal conductivity of solids. At the Debye temperature, the shortest wavelength of phonon corresponds to the size of the unit cell for simple cubic crystals.\(^{35}\)

At a temperature above the Debye temperature, the heat capacity and the wave velocity are both approximate constant, whilst the mean free path for crystalline dielectric materials is nearly proportional to \( 1/T \).\(^{36}\) In this case, the thermal conductivities of \( Y_2SiO_5 \) can be fitted perfectly by a \( 1/T \) equation in the range over 300–973 K. For the temperature above 973 K, the thermal conductivity should be a plateau and become independent of temperature. However, the actual condition is that the thermal conductivity reached a higher value 1.50 W/m K after a small platform in 873–1073 K. The abnormal increasing of thermal conductivity at temperatures higher than 1073 K may be resulted by the non-blocked radioactive heat transport through a material during measurement.\(^{37}\) The measured minimum thermal conductivity value is 1.34 W/m K and this value sustained in a temperature range from 873 to 1073 K.

For most materials, such a minimum plateau value of thermal conductivity should be the minimum thermal conductivity \( \kappa_{\text{min}} \). In selecting candidate materials for coating, it is necessary to know what the plateau value is and what temperature the transition from the \( 1/T \)-dependance to plateau behavior occurs.\(^{37}\) But experimentally the actual \( \kappa_{\text{min}} \) is very sensitive to the concentration of defects and preparation procedures, which make the comparison of thermal conductivity and the selection of material become difficult. Here, a calculated minimum thermal conductivity was used as a standard to choose suitable material for ETBC application.\(^{5,37}\)

The Debye temperature, \( \theta_D \), of a solid is always associated with its melting point and can be calculated from the average sound velocity, \( v_m \),\(^{35}\):

\[
\theta_D = \frac{h}{k_B} \left[ \frac{3n}{4\pi} \left( \frac{N\Lambda d}{M} \right)^{1/3} \right]^{1/3} v_m
\]  
(7)

The average sound velocity is defined as\(^{37}\):

\[
v_m = A \sqrt{\frac{E}{d}}
\]  
(8)

where \( M \) is the molecular mass, \( n \) the number of atoms in molecular, \( d \) the theoretical density, and \( E \) is Young’s Modulus. \( A \) has a value of 0.87 ± 0.02. At the Debye temperature, the shortest wavelength of phonon is corresponding to the size of the primitive cell, \( a_0 \).

\[
\theta_D = \frac{2v_m}{a_0 k_B}
\]  
(9)

The Debye temperature of \( Y_2SiO_5 \) can be calculated by Eqs. (7) and (8), and the value is 580 K. At the temperature higher than \( \theta_D \), the thermal conductivity approaches a minimum which is given by:

\[
\kappa \to \kappa_B v_m \Lambda \to \kappa_B v_m \Lambda_{\text{min}}
\]  
(10)

where \( \Lambda_{\text{min}} \) is the minimum free path, which equals to the average atomic spacing. Ignoring the fact that the material of interest contain more than one type of atom per unit cell or per formula unit, an equivalent atom having a mean atomic mass given by \( M/n \), where \( M \) is the molecular mass and \( n \) is the number of atoms per molecular, was employed to replace the different atoms in a molecule.\(^{37}\)

Combining Eqs. (7)–(10), the minimum thermal conductivity can be expressed as\(^{5,37}\):

\[
\kappa_{\text{min}} \to 0.87k_B N_{A}^{2/3} \pi^{1/3} E^{1/2} \frac{1}{M^{1/3}d^{1/6}}
\]  
(11)

The calculated minimum thermal conductivities of silicium carbide and selected EBC/TBC candidates are listed in Table 2,\(^{5,32–34,37}\) wherein \( Y_2SiO_5 \) shows a low value of 1.13 W/m K, being close to the measured minimum value of 1.34 W/m K.

### 3.3. Potential coating applications for \( Y_2SiO_5 \)

The low thermal conductivity of \( Y_2SiO_5 \) makes it very competitive as EBC/TBC candidates, especially when considering the match of thermal expansion coefficients between coating and substrate. At present, nickel-based superalloy is the most commonly used substrate material for gas turbine engines. Materials for the TBC layer are desired to have a thermal expansion coefficient close to that of nickel-based superalloys and have the lowest possible \( k \) value.

Fig. 7 presents the thermal expansion coefficient as a function of thermal conductivity for some coating materials and...
nickel-based superalloys. The nickel-based alloys used for turbine blades have a TEC from 14.0 to 16.0 × 10⁻⁶ K⁻¹. In Fig. 7, YSZ has a TEC of ∼9.0 × 10⁻⁶ K⁻¹ and Y₂SiO₅ of ∼8.4 × 10⁻⁶ K⁻¹. A very similar thermal expansion coefficient is noticed between Y₂SiO₅ and nickel-based superalloy, guaranteeing a low thermal stress at the coating and substrate interfaces and thereby a good thermal cycling resistance. According to Fig. 7, Y₂SiO₅ can be regarded as the most suitable coating material for nickel-based superalloys if only considering the thermal expansion and thermal conductivity. Moreover, combining the low thermal conductivity, thermal expansion coefficient and good corrosion resistance, Y₂SiO₅ is also an excellent environmental barrier coating on YSZ TBC layer. As a source of Y₃⁺, Y₂SiO₅ EBC can prolong the lifetime of zirconia-based ceramics with potential high-temperature structural applications. As shown in Table 3, Y₂SiO₅ in combination with Y₂Si₂O₇ sometimes.

### Table 3

<table>
<thead>
<tr>
<th>Compound</th>
<th>CTE</th>
<th>Compound</th>
<th>CTE</th>
<th>Compound</th>
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<tr>
<td>CrN</td>
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<td>β-SiC</td>
<td>5.1</td>
<td>TiC</td>
<td>8.3</td>
</tr>
<tr>
<td>Si₃N₄</td>
<td>3.7</td>
<td>B₄C</td>
<td>5.5</td>
<td>Y₂SiO₅</td>
<td>8.36</td>
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<tr>
<td>γ-Y₂Si₂O₇</td>
<td>3.90</td>
<td>ZrC</td>
<td>6.6</td>
<td>Mo₂C</td>
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<tr>
<td>α-SiC</td>
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<td>7.7</td>
<td>TiₓSiₓC</td>
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<tr>
<td>WC</td>
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<td>MoSi₂</td>
<td>8.2</td>
<td>TiₓAlₓC</td>
<td>9.0</td>
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</tbody>
</table>

In summary, the thermal properties of Y₂SiO₅ were systematically studied, including thermal expansion, thermal diffusivity, heat capacity and thermal conductivity. The linear thermal expansion coefficient of the single-phase polycrystalline bulk Y₂SiO₅ material determined by the push-rod dilatometer is (8.36 ± 0.5) × 10⁻⁶ K⁻¹ in the temperature range from 300 to 1573 K. During the test, no polymorph transformation was detected either in the heating or the cooling process.

The measured minimum thermal conductivity of Y₂SiO₅ is 1.34 W/m K at 873 K to 1173 K, which is lower than most commonly used thermal barrier coating material such as YSZ. The low thermal conductivity, low oxygen permeability as well as a relatively high thermal expansion coefficient ensure Y₂SiO₅ being a very competitive candidate material for both environmental/thermal barrier and oxygen-resistant coatings, especially in combination with Y₂Si₂O₇ sometimes.

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